

Equations of two-dimensional flow and bed deformation in general coordinate system

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Basic Equations of 2D Flow in (x, y) Co-orthogonal Coordinate System

$$\frac{\partial h}{\partial t} + \frac{\partial(hu)}{\partial x} + \frac{\partial(hv)}{\partial y} = 0 \quad (1)$$

$$\frac{\partial(uh)}{\partial t} + \frac{\partial(hu^2)}{\partial x} + \frac{\partial(huv)}{\partial y} = -hg \frac{\partial H}{\partial x} - \frac{\tau_x}{\rho} + D^x \quad (2)$$

$$\frac{\partial(vh)}{\partial t} + \frac{\partial(huv)}{\partial x} + \frac{\partial(hv^2)}{\partial y} = -hg \frac{\partial H}{\partial y} - \frac{\tau_y}{\rho} + D^y \quad (3)$$

in which,

$$\frac{\tau_x}{\rho} = C_d u \sqrt{u^2 + v^2} \quad \frac{\tau_y}{\rho} = C_d v \sqrt{u^2 + v^2} \quad (4)$$

$$D^x = \frac{\partial}{\partial x} \left[\nu_t \frac{\partial(uh)}{\partial x} \right] + \frac{\partial}{\partial y} \left[\nu_t \frac{\partial(uh)}{\partial y} \right] \quad (5)$$

$$D^y = \frac{\partial}{\partial x} \left[\nu_t \frac{\partial(vh)}{\partial x} \right] + \frac{\partial}{\partial y} \left[\nu_t \frac{\partial(vh)}{\partial y} \right] \quad (6)$$



Flow Equations in General Coordinate System

$$\frac{\partial}{\partial t} \left(\frac{h}{J} \right) + \frac{\partial}{\partial \xi} \left(\frac{hu^\xi}{J} \right) + \frac{\partial}{\partial \eta} \left(\frac{hu^\eta}{J} \right) = 0 \quad (7)$$

$$\begin{aligned} \frac{\partial u^\xi}{\partial t} + u^\xi \frac{\partial u^\xi}{\partial \xi} + u^\eta \frac{\partial u^\xi}{\partial \eta} + \alpha_1 u^\xi u^\xi + \alpha_2 u^\xi u^\eta + \alpha_3 u^\eta u^\eta = \\ -g \left[(\xi_x^2 + \xi_y^2) \frac{\partial H}{\partial \xi} + (\xi_x \eta_x + \xi_y \eta_y) \frac{\partial H}{\partial \eta} \right] \\ - \frac{C_d u^\xi}{hJ} \sqrt{(\eta_y u^\xi - \xi_y u^\eta)^2 + (-\eta_x u^\xi + \xi_x u^\eta)^2} + D^\xi \end{aligned} \quad (8)$$

$$\begin{aligned} \frac{\partial u^\eta}{\partial t} + u^\xi \frac{\partial u^\eta}{\partial \xi} + u^\eta \frac{\partial u^\eta}{\partial \eta} + \alpha_4 u^\xi u^\xi + \alpha_5 u^\xi u^\eta + \alpha_6 u^\eta u^\eta = \\ -g \left[(\eta_x \xi_x + \eta_y \xi_y) \frac{\partial H}{\partial \xi} + (\eta_x^2 + \eta_y^2) \frac{\partial H}{\partial \eta} \right] \\ - \frac{C_d u^\eta}{hJ} \sqrt{(\eta_y u^\xi - \xi_y u^\eta)^2 + (-\eta_x u^\xi + \xi_x u^\eta)^2} + D^\eta \end{aligned} \quad (9)$$

in which,

$$\alpha_1 = \xi_x \frac{\partial^2 x}{\partial \xi^2} + \xi_y \frac{\partial^2 y}{\partial \xi^2}, \quad \alpha_2 = 2 \left(\xi_x \frac{\partial^2 x}{\partial \xi \partial \eta} + \xi_y \frac{\partial^2 y}{\partial \xi \partial \eta} \right), \quad \alpha_3 = \xi_x \frac{\partial^2 x}{\partial \eta^2} + \xi_y \frac{\partial^2 y}{\partial \eta^2} \quad (10)$$

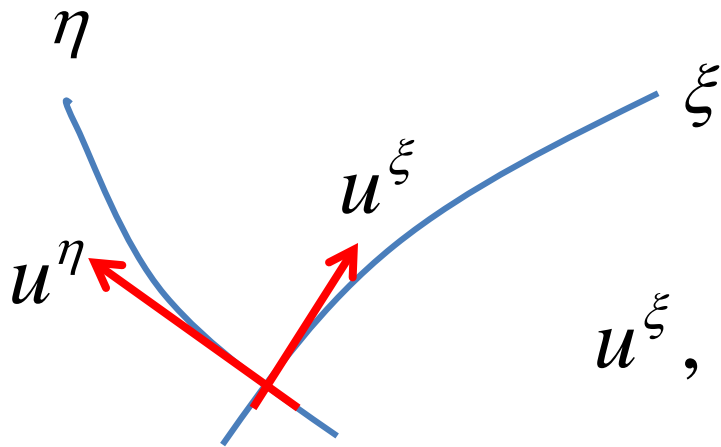
$$\alpha_4 = \eta_x \frac{\partial^2 x}{\partial \xi^2} + \eta_y \frac{\partial^2 y}{\partial \xi^2}, \quad \alpha_5 = 2 \left(\eta_x \frac{\partial^2 x}{\partial \xi \partial \eta} + \eta_y \frac{\partial^2 y}{\partial \xi \partial \eta} \right), \quad \alpha_6 = \eta_x \frac{\partial^2 x}{\partial \eta^2} + \eta_y \frac{\partial^2 y}{\partial \eta^2} \quad (11)$$

$$D^\xi =$$

$$\left(\xi_x \frac{\partial}{\partial \xi} + \eta_x \frac{\partial}{\partial \eta} \right) \left[\nu_t \left(\xi_x \frac{\partial u^\xi}{\partial \xi} + \eta_x \frac{\partial u^\xi}{\partial \eta} \right) \right] + \left(\xi_y \frac{\partial}{\partial \xi} + \eta_y \frac{\partial}{\partial \eta} \right) \left[\nu_t \left(\xi_y \frac{\partial u^\xi}{\partial \xi} + \eta_y \frac{\partial u^\xi}{\partial \eta} \right) \right] \quad (12)$$

$$D^\eta =$$

$$\left(\xi_x \frac{\partial}{\partial \xi} + \eta_x \frac{\partial}{\partial \eta} \right) \left[\nu_t \left(\xi_x \frac{\partial u^\eta}{\partial \xi} + \eta_x \frac{\partial u^\eta}{\partial \eta} \right) \right] + \left(\xi_y \frac{\partial}{\partial \xi} + \eta_y \frac{\partial}{\partial \eta} \right) \left[\nu_t \left(\xi_y \frac{\partial u^\eta}{\partial \xi} + \eta_y \frac{\partial u^\eta}{\partial \eta} \right) \right] \quad (13)$$



u^ξ, u^η ; contra-variant velocity components

Bed shear stress

Depth averaged total velocity

$$\begin{pmatrix} u \\ v \end{pmatrix} = \frac{1}{J} \begin{pmatrix} \eta_y & -\xi_y \\ -\eta_x & \xi_x \end{pmatrix} \begin{pmatrix} u^\xi \\ u^\eta \end{pmatrix} \longrightarrow V = \sqrt{u^2 + v^2}$$

$$V = \frac{1}{n_m} h^{2/3} I_e^{1/2} \longrightarrow I_e = \frac{V^2 n_m^2}{h^{4/3}} \quad \left[\begin{array}{l} u_* = \sqrt{ghI_e} \\ \text{or} \\ \tau_b = \rho ghI_e \end{array} \right.$$

or

$$V = C \sqrt{hI_e} \longrightarrow I_e = \frac{V^2}{C^2 h}$$

$$\tau_* = \frac{\tau_b}{\rho s g d} \quad \text{or} \quad \tau_* = \frac{u_*^2}{s g d} \quad \text{or} \quad \tau_* = \frac{h I_e}{s d}$$

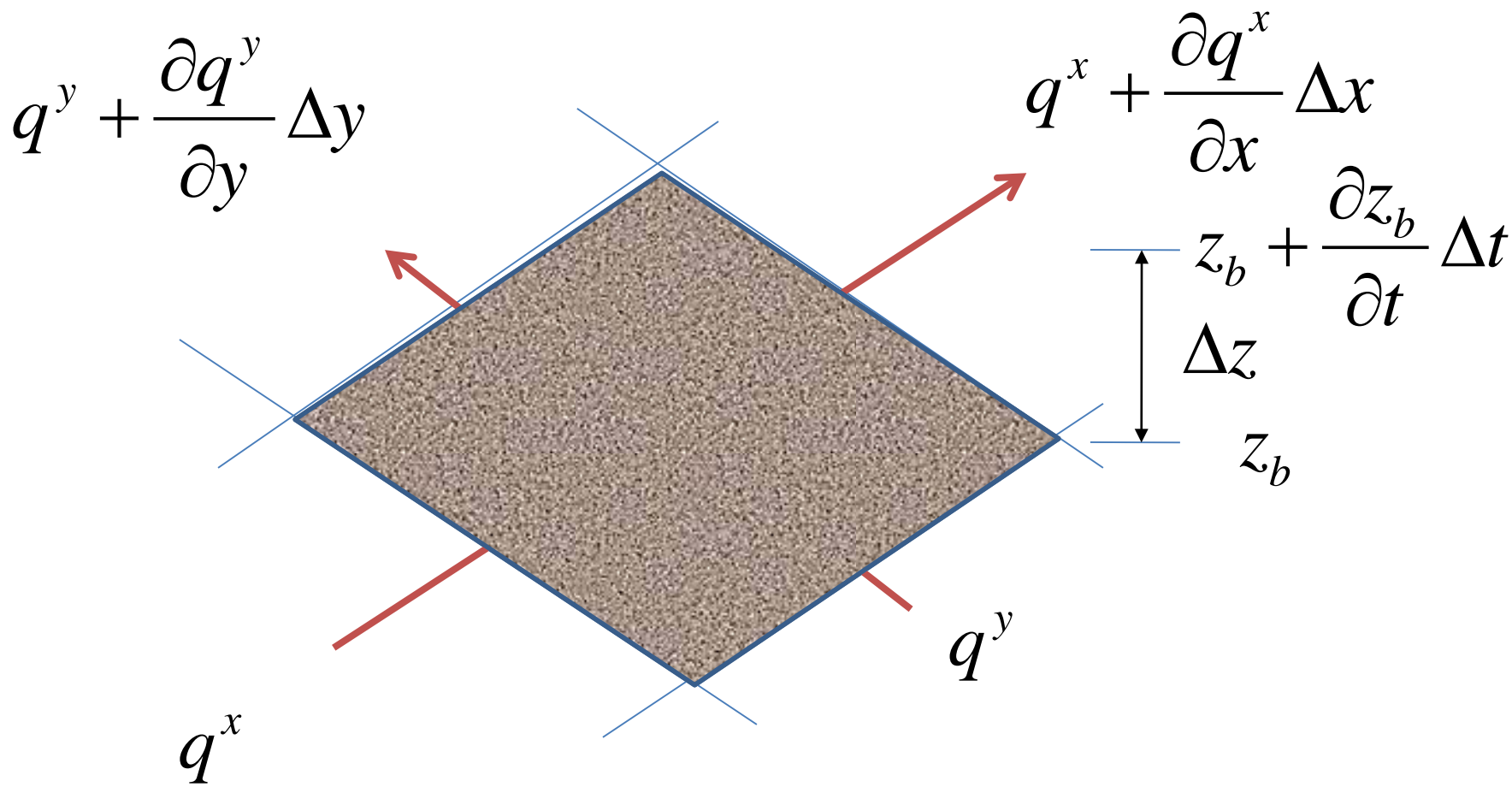
In which $s = 1 - \frac{\rho_s}{\rho}$

Submerged relative weight of sand
(=1.65)

Bed load transport rate in depth averaged flow direction

$$q_b = 17\tau_*^{3/2} \left(1 - \frac{\tau_{*c}}{\tau_*}\right) \left(1 - \sqrt{\frac{\tau_{*c}}{\tau_*}}\right) \sqrt{sgd^3}$$

Ashida, K. and M. Michiue, 1972, Study on hydraulic resistance and bedload transport rate in alluvial streams, *Transactions*, Japan Society of Civil Engineering, 206: 59-69 (in Japanese)



$$\frac{\partial z_b}{\partial t} + \frac{1}{1 - \lambda} \left[\frac{\partial q^x}{\partial x} + \frac{\partial q^y}{\partial y} \right] = 0$$

2-dimensional Continuity Equations for Bedload

$$\frac{\partial z_b}{\partial t} + \frac{1}{1 - \lambda} \left[\frac{\partial q^x}{\partial x} + \frac{\partial q^y}{\partial y} \right] = 0 \quad (25)$$

in which, z_b is bed elevation, q^x and q^y are bedload transport rate per unit width in x and y directions, and λ is void ratio of bed material.

$$\frac{\partial}{\partial t} \left(\frac{z_b}{J} \right) + \frac{1}{1 - \lambda} \left[\frac{\partial}{\partial \xi} \left(\frac{q^\xi}{J} \right) + \frac{\partial}{\partial \eta} \left(\frac{q^\eta}{J} \right) \right] = 0 \quad (26)$$

in which q^ξ and q^η are contravariant components of bedload sediment transport rate in ξ and η direction. They are also needed to be transformed as follows to describe in actual sediment transport rate in [Length²/Time].

$$\widetilde{q}^\xi = \frac{q^\xi}{\xi_r}, \quad \widetilde{q}^\eta = \frac{q^\eta}{\eta_r} \quad (27)$$

Watanabe et al.[2] proposed the following equation considering the gravitational effect in streamline and transverse directions.

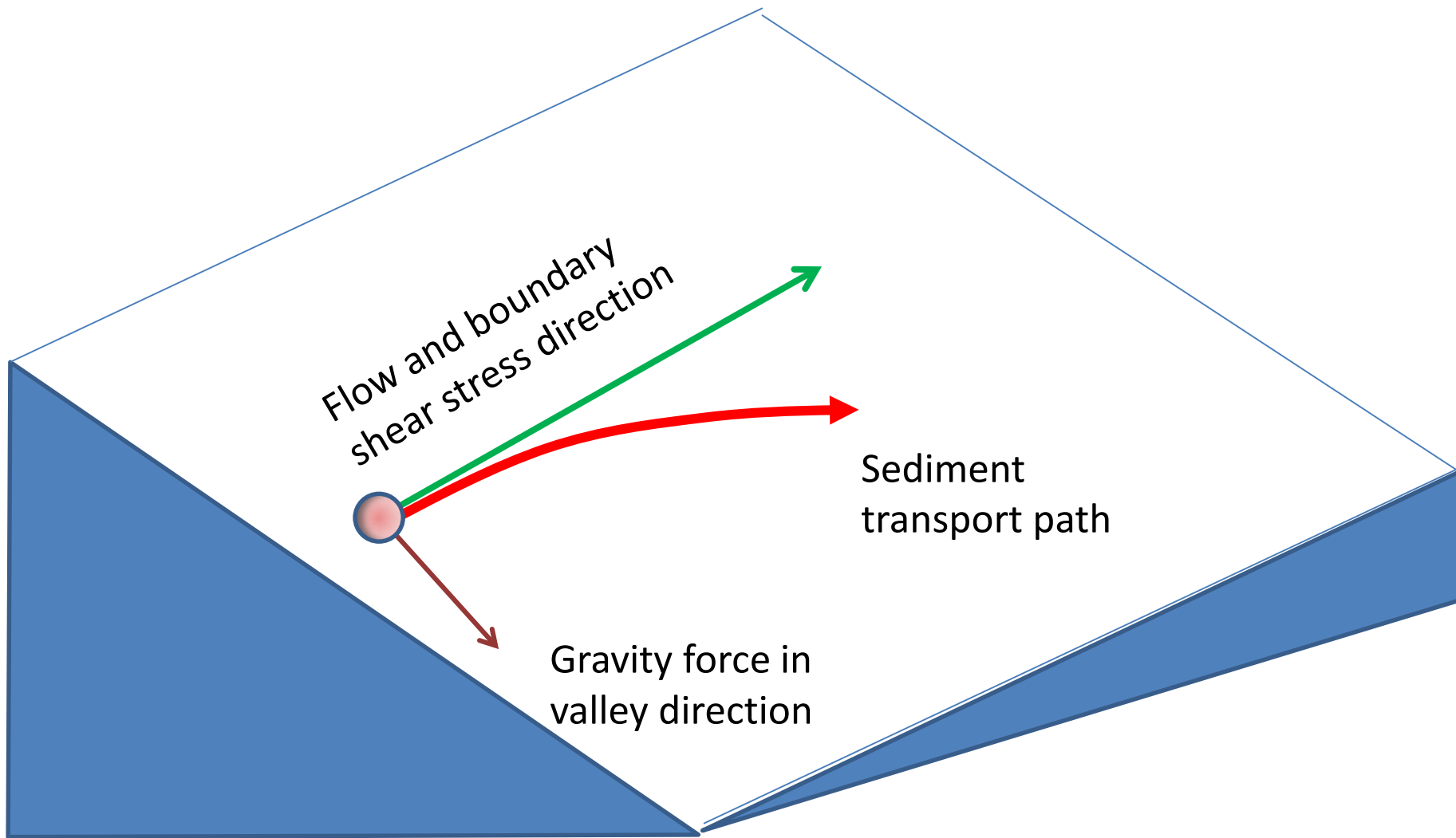
$$\tilde{q}^{\xi} = q_b \left[\frac{\tilde{u}_b^{\xi}}{V_b} - \gamma \left(\frac{\partial z_b}{\partial \tilde{\xi}} + \cos \theta \frac{\partial z_b}{\partial \tilde{\eta}} \right) \right] \quad (32)$$

$$\tilde{q}^{\eta} = q_b \left[\frac{\tilde{u}_b^{\eta}}{V_b} - \gamma \left(\frac{\partial z_b}{\partial \tilde{\eta}} + \cos \theta \frac{\partial z_b}{\partial \tilde{\xi}} \right) \right] \quad (33)$$

in which, \tilde{u}_b^{ξ} and \tilde{u}_b^{η} are the velocity components at the bottom in ξ and η directions, V_b is the total velocity at the bottom, θ is an angle between ξ -axis and η -axis. γ is an adjustment coefficient for slope gravitational effect. Hasegawa[3] proposed the following formula.

$$\gamma = \sqrt{\frac{\tau_{*c}}{\mu_s \mu_k \tau_*}} \quad (34)$$

in which, μ_s and μ_k are static and kinetic friction coefficient of bed material.



Flow and boundary
shear stress direction

Sediment
transport path

Gravity force in
valley direction

Watanabe et al.[2] proposed the following equation considering the gravitational effect in streamline and transverse directions.

$$\tilde{q}^{\xi} = q_b \left[\frac{\tilde{u}_b^{\xi}}{V_b} - \gamma \left(\frac{\partial z_b}{\partial \tilde{\xi}} + \cos \theta \frac{\partial z_b}{\partial \tilde{\eta}} \right) \right] \quad (32)$$

$$\tilde{q}^{\eta} = q_b \left[\frac{\tilde{u}_b^{\eta}}{V_b} - \gamma \left(\frac{\partial z_b}{\partial \tilde{\eta}} + \cos \theta \frac{\partial z_b}{\partial \tilde{\xi}} \right) \right] \quad (33)$$

in which, \tilde{u}_b^{ξ} and \tilde{u}_b^{η} are the velocity components at the bottom in ξ and η directions, V_b is the total velocity at the bottom, θ is an angle between ξ -axis and η -axis. γ is an adjustment coefficient for slope gravitational effect.

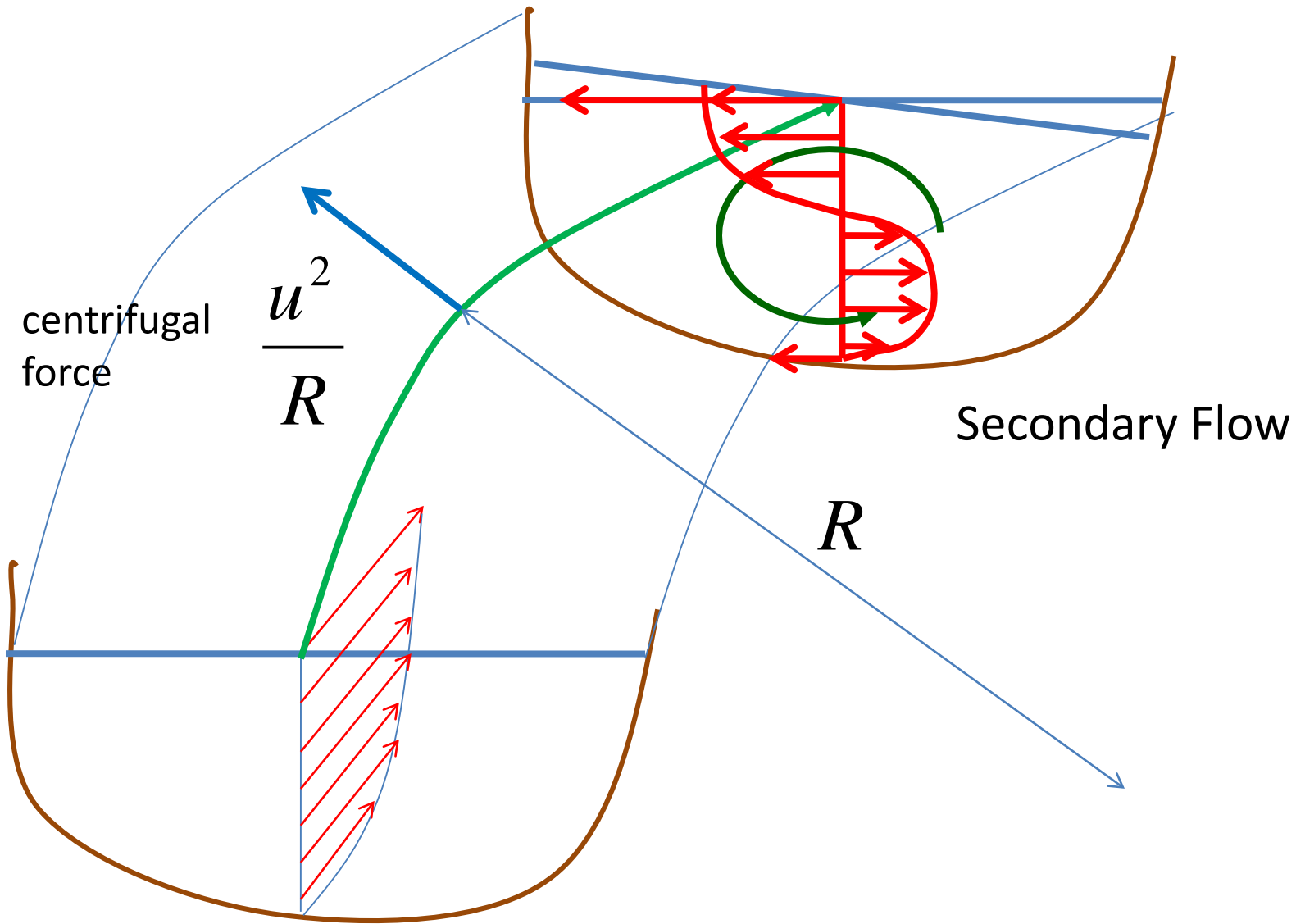
Watanabe et al., Groin Arrangements made of natural willows for reducing bed deformation in a curved channel., Advance in River Engineering, Vol.7, pp.285-290, 2001, JSCE (in Japanese).

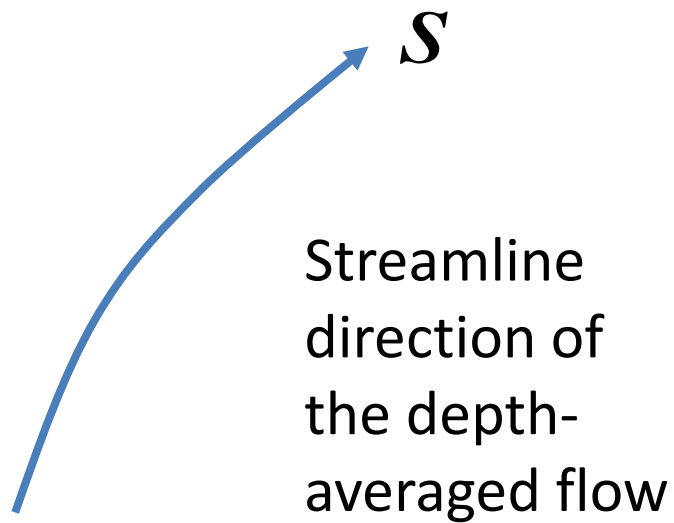
between ξ -axis and η -axis. γ is an adjustment coefficient for slope gravitational effect. Hasegawa[3] proposed the following formula.

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in which, μ_s and μ_k are static and kinetic friction coefficient of bed material.

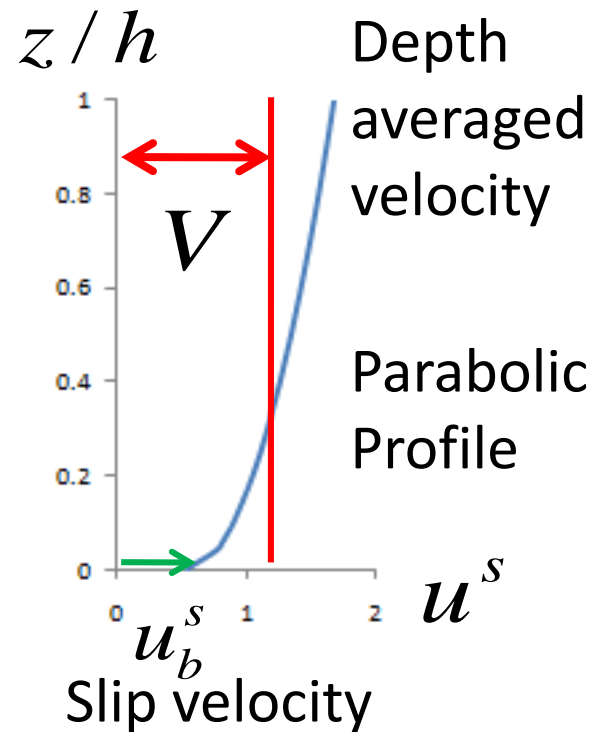
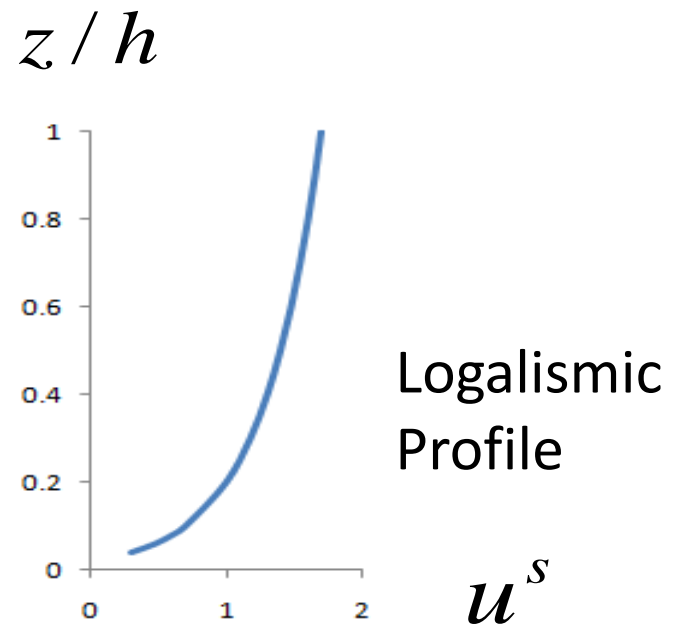
Hasegawa, K. and Yamaoka, S.: The Effect of plane and bed forms of channels upon the meander development, Journal of Hydraulic, Coastal and Environmental Engineering, JSCE, Vol229, pp.143--152, 1980. (in Japanese)





$$u_b^s = \beta V$$

Engelund (1974)



Velocity components at channel bottom

The following simple relation is assumed between depth averaged flow velocities and bottom velocities.

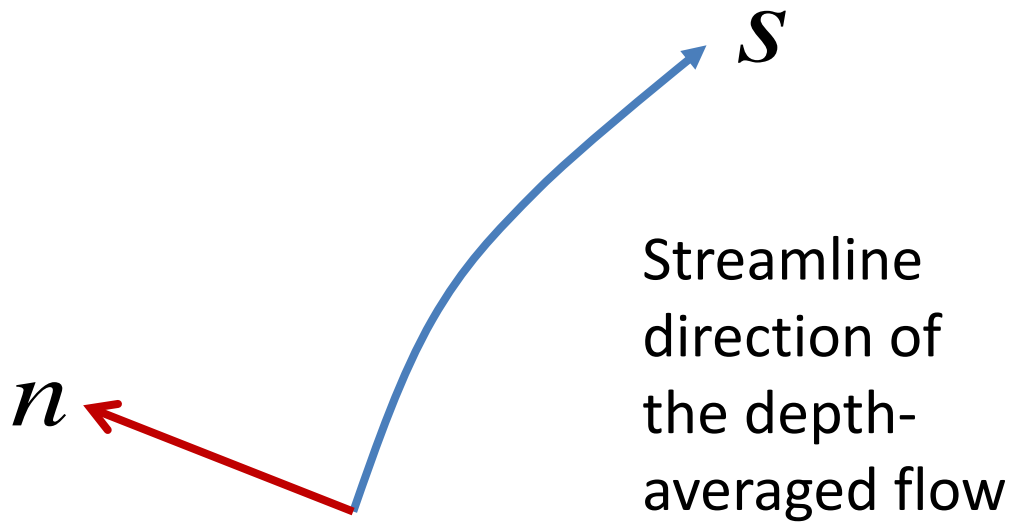
$$\tilde{u}_b^s = \beta V \quad (35)$$

in which, \tilde{u}_b^s is bottom velocity along the depth averaged stream line. Engelund[4] used a parabolic function for velocity profile in depth direction, and proposed the following function.

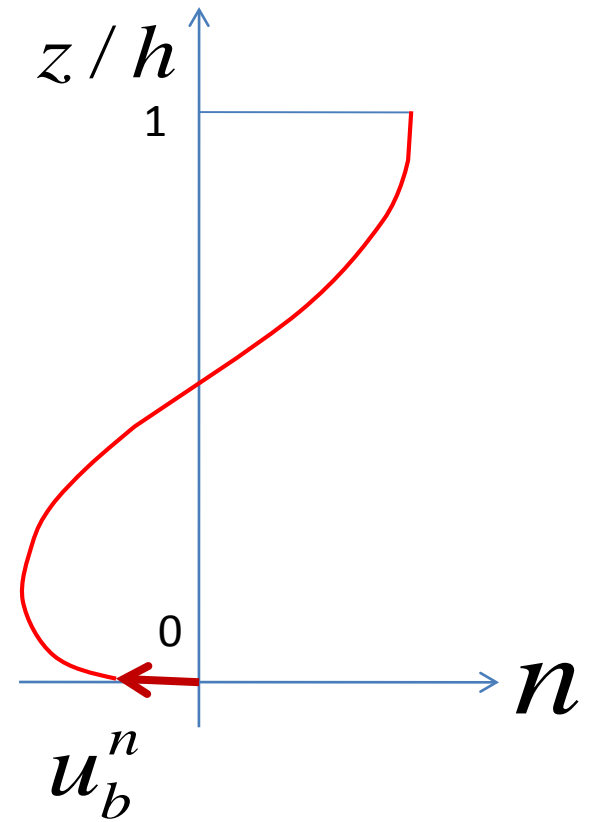
$$\beta = 3(1 - \sigma)(3 - \sigma), \quad \sigma = \frac{3}{\phi_0 \kappa + 1} \quad (36)$$

in which, ϕ_0 is velocity coefficient(= V/u_*), κ Von Karman's constant(=0.4).

Engelund, F.: Flow and Bed Topography in Channel Bend, Jour. of Hydr. Div., ASCE, Vol.100. HY11, pp.1631--1648, 1974.



Normal direction
to the streamline



$$u_b^n = u_b^s N_* \frac{h}{r_s}$$

Emgelund (1974)

When the stream line is curved, the secondary flow, or spiral flow is generated. The following equation is used to estimate the velocity components considering secondarily flow.

$$\widetilde{u}_b^n = \widetilde{u}_b^s N_* \frac{h}{r_s} \quad (37)$$

in which, \widetilde{u}_b^n is a bottom velocity perpendicular to the direction of stream line, which is positive 90 degree clock wise direction from the stream line direction, r_s is a radius of curvature of the streamline, N_* is a constant (=7, Engelund[4]).

Intensity of the secondary flow is proportional to the depth and inverse proportional to the radius of the curvature of the streamline.